Chapter 6 Oscillations

The potential energy then reduces to

$$V = \frac{k}{2}(\eta_2 - \eta_1)^2 + \frac{k}{2}(\eta_3 - \eta_2)^2,$$

OF

$$V = \frac{k}{2}(\eta_1^2 + 2\eta_2^2 + \eta_3^2 - 2\eta_1\eta_2 - 2\eta_2\eta_3). \tag{6.49}$$

Hence, the V tensor has the form

$$\mathbf{V} = \begin{pmatrix} k & -k & 0 \\ -k & 2k & -k \\ 0 & -k & k \end{pmatrix}. \tag{6.50}$$

The kinetic energy has an even simpler form:

$$T = \frac{m}{2}(\dot{\eta}_1^2 + \dot{\eta}_3^2) + \frac{M}{2}\dot{\eta}_2^2,\tag{6.51}$$

so that the T tensor is diagonal:

$$\mathbf{T} = \begin{pmatrix} m & 0 & 0 \\ 0 & M & 0 \\ 0 & 0 & m \end{pmatrix}. \tag{6.52}$$

Combining these two tensors, the secular equation appears as

$$|\mathbf{V} - \omega^2 \mathbf{T}| = \begin{vmatrix} k - \omega^2 m & -k & 0 \\ -k & 2k - \omega^2 M & -k \\ 0 & -k & k - \omega^2 m \end{vmatrix} = 0.$$
 (6.53)

Direct evaluation of the determinant leads to the cubic equation in ω^2 :

$$\omega^{2}(k - \omega^{2}m)(k(M + 2m) - \omega^{2}Mm) = 0, \qquad (6.54)$$

with the obvious solutions

$$\omega_1 = 0$$
, $\omega_2 = \sqrt{\frac{k}{m}}$, $\omega^3 = \sqrt{\frac{k}{m} \left(1 + \frac{2m}{M}\right)}$. (6.55)

The first eigenvalue, $\omega_1 = 0$, may appear somewhat surprising and even alarming at first sight. Such a solution does not correspond to an oscillatory motion at all, for the equation of motion for the corresponding normal coordinate is

$$\ddot{\zeta}_1 = 0$$
,

which produces a uniform translational motion. But this is precisely the key to the difficulty. The vanishing frequency arises from the fact that the molecule

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Julia Schneider

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Table of Contents Classical Mechanics Chapter 6 Homework Solutions

- 1. Understanding the eBook Classical Mechanics Chapter 6 Homework Solutions
 - The Rise of Digital Reading Classical Mechanics Chapter 6 Homework Solutions
 - Advantages of eBooks Over Traditional Books
- 2. Identifying Classical Mechanics Chapter 6 Homework Solutions
 - Exploring Different Genres
 - Considering Fiction vs. Non-Fiction
 - Determining Your Reading Goals
- 3. Choosing the Right eBook Platform
 - Popular eBook Platforms
 - Features to Look for in an Classical Mechanics Chapter 6 Homework Solutions
 - User-Friendly Interface
- 4. Exploring eBook Recommendations from Classical Mechanics Chapter 6 Homework Solutions
 - Personalized Recommendations
 - Classical Mechanics Chapter 6 Homework Solutions User Reviews and Ratings
 - Classical Mechanics Chapter 6 Homework Solutions and Bestseller Lists
- 5. Accessing Classical Mechanics Chapter 6 Homework Solutions Free and Paid eBooks
 - Classical Mechanics Chapter 6 Homework Solutions Public Domain eBooks
 - Classical Mechanics Chapter 6 Homework Solutions eBook Subscription Services
 - Classical Mechanics Chapter 6 Homework Solutions Budget-Friendly Options

- 6. Navigating Classical Mechanics Chapter 6 Homework Solutions eBook Formats
 - o ePub, PDF, MOBI, and More
 - Classical Mechanics Chapter 6 Homework Solutions Compatibility with Devices
 - Classical Mechanics Chapter 6 Homework Solutions Enhanced eBook Features
- 7. Enhancing Your Reading Experience
 - o Adjustable Fonts and Text Sizes of Classical Mechanics Chapter 6 Homework Solutions
 - Highlighting and Note-Taking Classical Mechanics Chapter 6 Homework Solutions
 - Interactive Elements Classical Mechanics Chapter 6 Homework Solutions
- 8. Staying Engaged with Classical Mechanics Chapter 6 Homework Solutions
 - Joining Online Reading Communities
 - Participating in Virtual Book Clubs
 - Following Authors and Publishers Classical Mechanics Chapter 6 Homework Solutions
- 9. Balancing eBooks and Physical Books Classical Mechanics Chapter 6 Homework Solutions
 - Benefits of a Digital Library
 - Creating a Diverse Reading Collection Classical Mechanics Chapter 6 Homework Solutions
- 10. Overcoming Reading Challenges
 - Dealing with Digital Eye Strain
 - Minimizing Distractions
 - Managing Screen Time
- 11. Cultivating a Reading Routine Classical Mechanics Chapter 6 Homework Solutions
 - Setting Reading Goals Classical Mechanics Chapter 6 Homework Solutions
 - Carving Out Dedicated Reading Time
- 12. Sourcing Reliable Information of Classical Mechanics Chapter 6 Homework Solutions
 - Fact-Checking eBook Content of Classical Mechanics Chapter 6 Homework Solutions
 - Distinguishing Credible Sources
- 13. Promoting Lifelong Learning
 - Utilizing eBooks for Skill Development
 - Exploring Educational eBooks
- 14. Embracing eBook Trends
 - Integration of Multimedia Elements

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