#### Solutions to the problems in Chapter 6 and 7

#### 6.3 Pressure of a Fermi gas at zero temperature

The number of electrons N and the internal energy U, in the volume V, are

$$N = V \int_{0}^{\infty} D(\varepsilon) f(\varepsilon) d\varepsilon$$
,  $U = V \int_{0}^{\infty} \varepsilon D(\varepsilon) f(\varepsilon) d\varepsilon$ , (1)

The Fermi distribution function  $f(\varepsilon)$  and the density of states (per unit volume)  $D(\varepsilon)$  are

$$f(\varepsilon) = \frac{1}{e^{B(\varepsilon - \mu)} + 1}, \quad D(\varepsilon) = A\sqrt{\varepsilon}, \quad A = \frac{\sqrt{2nc^3}}{\pi^2 \hbar^3}.$$
 (2)

At zero temperature

$$f(\varepsilon) = \theta(\varepsilon_F - \varepsilon), \quad \varepsilon_F = \frac{\hbar^2 k_F^2}{2m}, \quad k_F = (3\pi^2 n)^{1/3}$$
 (3)

Introducing the step function and  $D(\varepsilon) = A\sqrt{\varepsilon}$  in (1), we may determine U(0) and an alternative expression for A:

$$N = V \int_0^{\varepsilon_F} A \sqrt{\varepsilon} d\varepsilon = \tfrac{2}{3} V A \varepsilon_F^{3/2} \ \Rightarrow \ A = \frac{3n}{2 \varepsilon_F^{3/2}}, \qquad U(0) = V \tfrac{2}{5} A \varepsilon_F^{5/2} = \tfrac{3}{5} N \varepsilon_F \ \ (4)$$

i.e.  $U \propto \varepsilon_{\rm F} \propto k_{\rm F}^2 \propto V^{-2/3}$  and the pressure at zero temperature is

$$P = -\frac{\partial U(0)}{\partial V} = \frac{2}{3} \frac{U}{V} = \frac{2}{5} n \varepsilon_{Y} = \frac{(3\pi^{2})^{2/3} n^{5/3} \hbar^{2}}{5m}$$
(5)

in sharp contrast to a classical gas, where the zero-temperature pressure is zero.

#### 6.4 Density of states in low dimensions

The density of states in  $\vec{k}$ -space is  $D_{\vec{k}}=2/(2\pi)^d$  in d dimensions (6.34)

$$D(\varepsilon) = \int [d\vec{k}] \delta(\varepsilon - \varepsilon_{\vec{k}}) \equiv \int D_{\vec{k}} \delta(\varepsilon - \varepsilon_{\vec{k}}) d\vec{k} = \frac{2}{(2\pi)^d} \int \delta(\varepsilon - \varepsilon_{\vec{k}}) d\vec{k}$$
 (1)

Introducing "spherical coordinates" in d dimensions and  $k = |\vec{k}|$ , then

$$d\vec{k} = \begin{cases} 2 dk & , & d = 1 \\ 2\pi k dk & , & d = 2 \\ 4\pi k^2 dk & , & d = 3 \end{cases}$$
(2)

The factor 2 in the one-dimensional case appears because the one-dimensional wave vector may assume both positive and negative values, whereas  $k \ge 0$  per definition.

$$\varepsilon_{\vec{k}} = \varepsilon_k = \frac{\hbar^2 k^2}{2m} \implies k = \frac{\sqrt{2m\varepsilon_{\vec{k}}}}{\hbar}, \quad dk = \left(\frac{m}{2\hbar^2\varepsilon_{\vec{k}}}\right)^{1/2} d\varepsilon_{\vec{k}} \implies (3)$$

$$D(\varepsilon) = \begin{cases} \frac{2}{2\pi} \int_0^\infty 2\,\delta(\varepsilon - \varepsilon_k) \left(\frac{m}{2h^2\varepsilon_k}\right)^{1/2} d\varepsilon_k = & \frac{\sqrt{2m}}{\pi\,h} \varepsilon^{-1/2}, \quad d = 1 \\ \frac{2}{(2\pi)^2} \int_0^\infty 2\pi \frac{\sqrt{2m}\varepsilon_k}{h} \,\delta(\varepsilon - \varepsilon_k) \left(\frac{m}{2h^2\varepsilon_k}\right)^{1/2} d\varepsilon_k = & \frac{m}{\pi\,h^2}, \qquad d = 2 \\ \frac{2}{(2\pi)^3} \int_0^\infty 4\pi \frac{2m\varepsilon_k}{h^2} \,\delta(\varepsilon - \varepsilon_k) \left(\frac{m}{2h^2\varepsilon_k}\right)^{1/2} d\varepsilon_k = & \frac{\sqrt{2m^3}}{\pi^2h^3} \varepsilon^{1/2}, \quad d = 3 \end{cases}$$

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**Amitabh Chandra** 

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