Classical Mechanics Problem 1: Central Potential Solution

a) Integrals of motion for a central potential V(r):

Angular Momentum $L = rv_t = r^2 \dot{\phi}$

Energy per unit mass $E = \frac{1}{2} (\hat{r}^2 + v_t^2) + V(r) = \frac{1}{2} \hat{r}^2 + V_{eff}(r)$

where v_r is the tangential velocity and V_{eff} is defined as

$$V_{\text{eff}}(r) = V(r) + \frac{L^2}{2r^2}$$

If the orbit is circular, the distance of the test body from the origin is invariant: $\dot{r} = 0$, which implies that the body is always at the equilibrium-distance:

$$\frac{dV_{\rm eff}}{dr} = 0$$
 \Rightarrow $\frac{dV}{dr} = \frac{L^2}{r^3} = \frac{v_1^3}{r} = r\dot{\phi}^2$

then

$$\dot{\phi} = \omega_{\phi} = \frac{L}{r^2} = \left(\frac{1}{r}\frac{dV}{dr}\right)^{1/2}$$

so for the period we get

$$T_{\phi} = \frac{2\pi}{\omega_{\phi}} = 2\pi \left(\frac{1}{r} \frac{dV}{dr}\right)^{-1/2}$$

b) Write the orbit as in the statement of the problem:

$$r(t) = r_0 + \epsilon(t)$$
 with $\frac{dV_{\rm eff}}{dr}(r_0) = 0$ and $\epsilon^2 \ll r_0^2$.

The energy per unit mass is now $E = \frac{1}{2}\tilde{\epsilon}^2 + V_{eff}(r_0 + \epsilon)$, and since ϵ is small we may Taylor-expand the potential as

$$V_{\text{eff}}(r_0 + \epsilon) = V_{\text{eff}}(r_0) + \underbrace{\frac{dV_{\text{eff}}}{dr}(r_0)\epsilon}_{=0} + \frac{1}{2} \frac{d^2V_{\text{eff}}}{dr^2}(r_0)\epsilon^2 + O(\epsilon^3)$$

see them.

$$E - V_{eff}(r_0) = \frac{1}{2}\dot{\epsilon}^2 + \frac{1}{2}\frac{d^2V_{eff}}{dr^2}(r_0)\epsilon^2 + O(\epsilon^3) = \text{const.}$$

In the above equation we readily recognize the equation of the simple harmonic oscillator with

$$\omega_r = \left(\frac{d^2V_{\text{eff}}}{dr^2}\right)_{r=r_0}^{1/2}$$

Classical Mechanics Problem 1 Central Potential Solution

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